

Direct-Drive Target Design: High Average Power Laser Program

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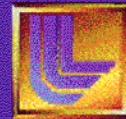


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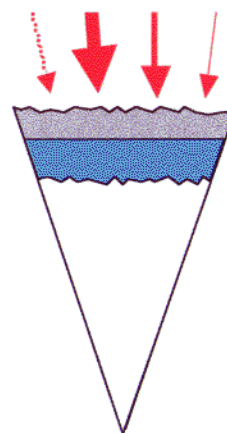
Critical Issues in the Design of A High Gain Direct-Drive Target for Inertial Fusion Energy



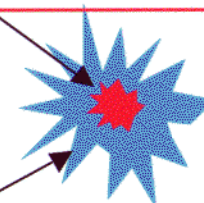
| | | |
|---|---|---|
| Laser-target coupling | ⇐ | Beam geom., intensity, λ , ablator design (scale lengths)... |
| Target isentrope control (ablator/fuel) | ⇐ | Pulse shaping... |
| Implosion symmetry (low modes) | ⇐ | Beam geometry and balance... |
| Implosion stability (higher modes) | ⇐ | Laser imprint, inner/outer surface finish, target build, adiabat control, λ , ... |
| Ignition and propagating burn | ⇐ | All of the above! |

High(er)-mode stability is (probably) the greatest source of uncertainty.

Success with high-gain direct-drive targets for IFE will depend fundamentally on our ability to achieve Rayleigh-Taylor stabilization



Ignition hotspot:
 $(\rho R.T)_{\text{hotspot}} \sim 0.3 \text{ g.cm}^{-2} @ 12 \text{ keV}$



Yield (gain) $\sim (\rho R)_{\text{fuel}} / [(\rho R)_{\text{fuel}} + 6]$
 $\Rightarrow (\rho R)_{\text{fuel}} \sim 2\text{-}2.5 \text{ g.cm}^{-2}$

It's the Wavelength Stupid!

DPSSLs (0.349 μm) –v– KrF (0.248 μm)



■ Higher ideal implosion velocity

$$V_{\text{implosion}} = V_{\text{exhaust}} \cdot \ln(m_f/m_i)$$

$$= P/(dm/dt) \cdot \ln(m_f/m_i) \sim (I \lambda^2)^{1/3} \cdot \ln(m_f/m_i)$$

■ Lower coupling (rocket) efficiency

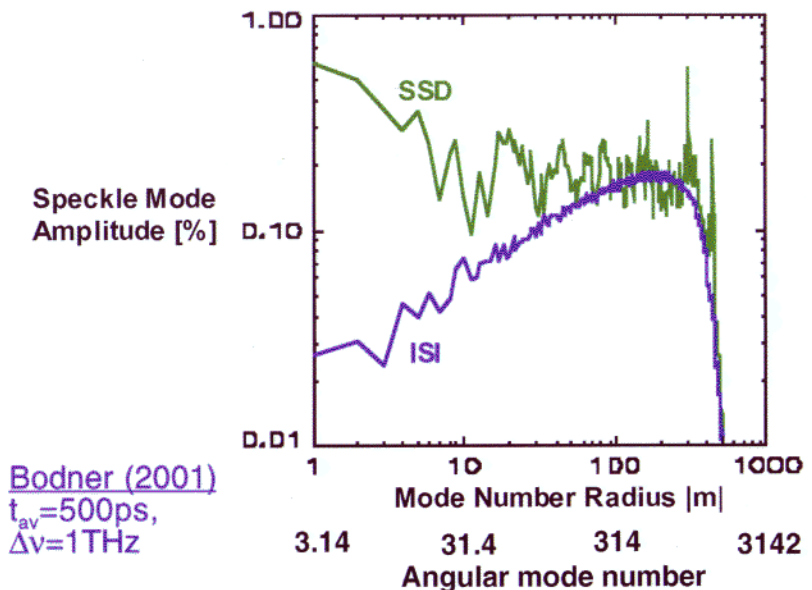
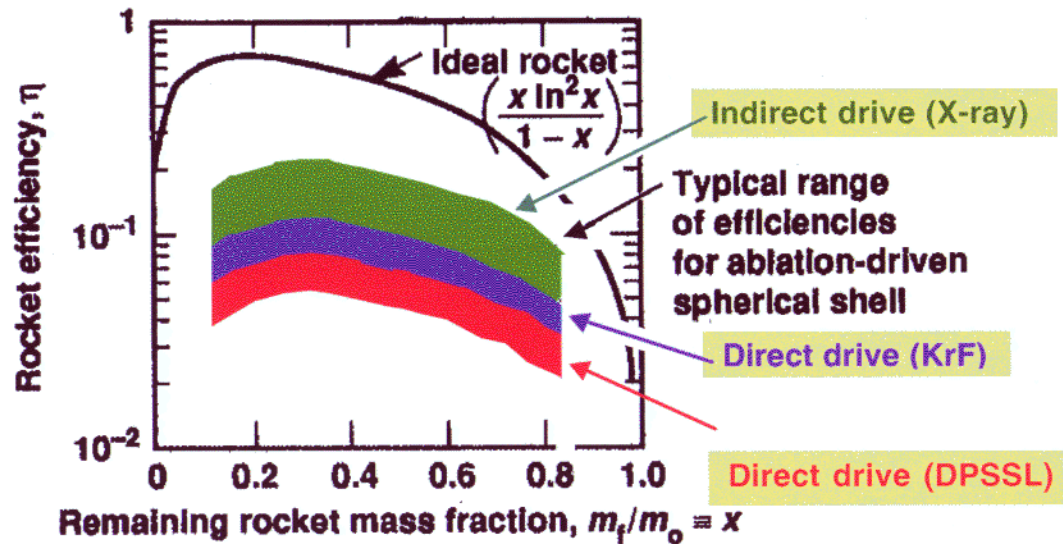
- Laser energy is deposited in target corona at $n \leq n_{\text{crit}} \sim 1/\lambda^2$ So factor of ~ 2 lower critical density for DPSSLs

■ Lower ablation velocity for stability

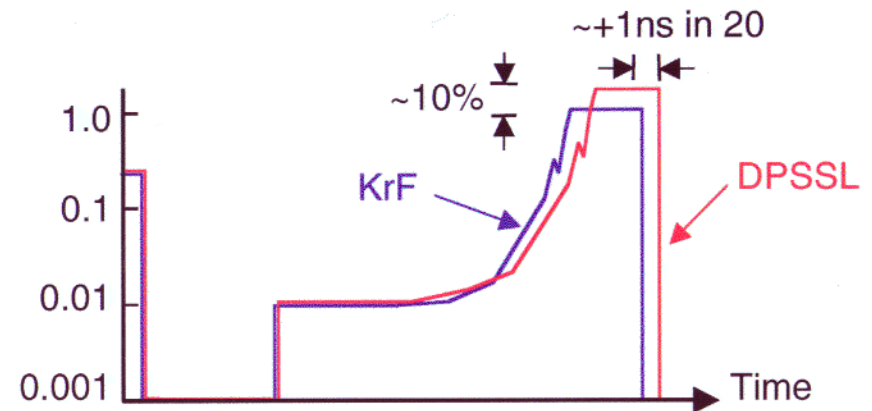
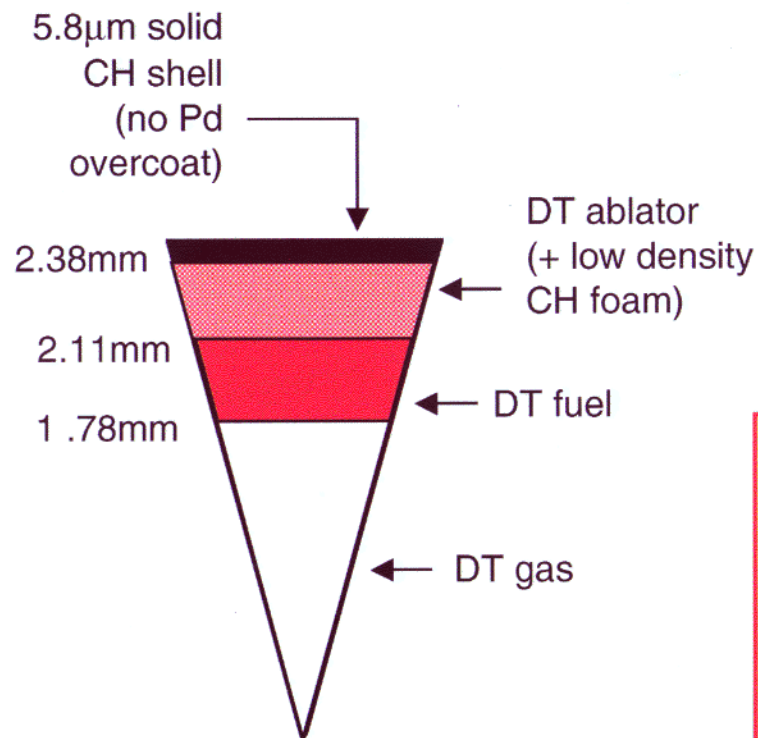
$$V_{\text{ablation}} = (dm/dt)/\rho \sim \alpha^{3/5} I^{-1/15} \lambda^{-14/15}$$

■ Less optimum (?) beam-smoothing

- DPSSLs use SSD
- KrF can use either ISI or SSD



KrF and DPSSL Give Comparable 1-D Performances within ~10%



| | <u>KrF</u> | <u>DPSSL</u> |
|----------------------|-----------------------------|------------------------|
| • λ | 0.248 μ m | 0.349 μ m |
| • Yield | 350MJ \longleftrightarrow | 350MJ |
| • Timing | 0 | ~+0.8ns |
| • E_{laser} | 2.9MJ | 3.2MJ |
| • Gain | 120 | 110 |
| • Max ρR | 2.11 g/cm ² | 2.18 g/cm ² |
| • KE margin | 31% | 29% |
| • Velocity | 3.30e7cm/s | 3.16e7cm/s |

Pulse Shape Tailoring of the Adiabatic Profile can Improve Stability without Compromising Ignition



- High gains and minimum drive energy for ignition require low fuel adiabats α (Herrmann et al), where: $\alpha = P/P_{\text{Fermi}}$



$$E_{\text{ign}} \sim \alpha^{1.8} V^{-6} P^{-0.8}$$

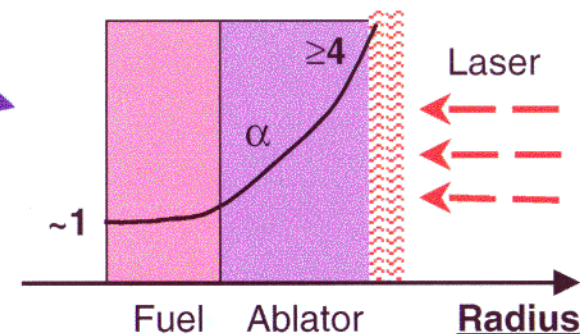
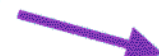
- But ablation velocity is a key factor in shell stability and increases with increasing adiabat



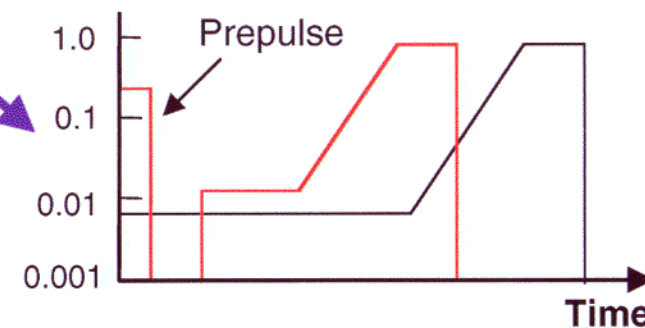
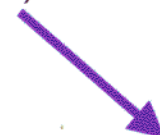
$$\gamma_{RT} \sim c_1 \sqrt{Akg} - c_2 kV_A$$

$$V_A \sim \frac{\dot{m}}{\rho} \sim \alpha^{3/5} I_{\text{laser}}^{-1/15}$$

- Tailoring the adiabat profile through the fuel and ablator can improve stability without compromising ignition and high gain



- Adiabat tailoring is achieved by a picket “stake” pulse – low energy prepulse with main pulse separated by power shut-off (Lindl, Verdon, Betti)

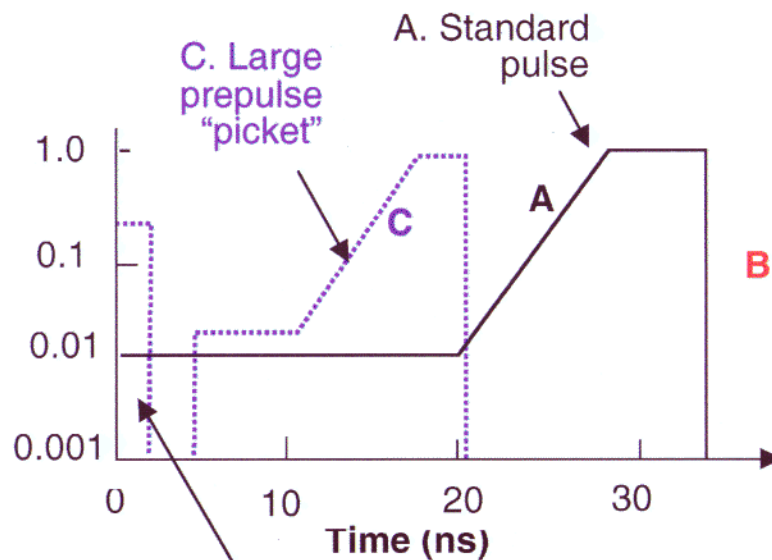


- Adiabat shaping may also be achievable with radiation preheat (Bodner)

Decaying Shock from Prepulse Produces High Ablator Adiabatic while Maintaining Low Fuel Adiabatic



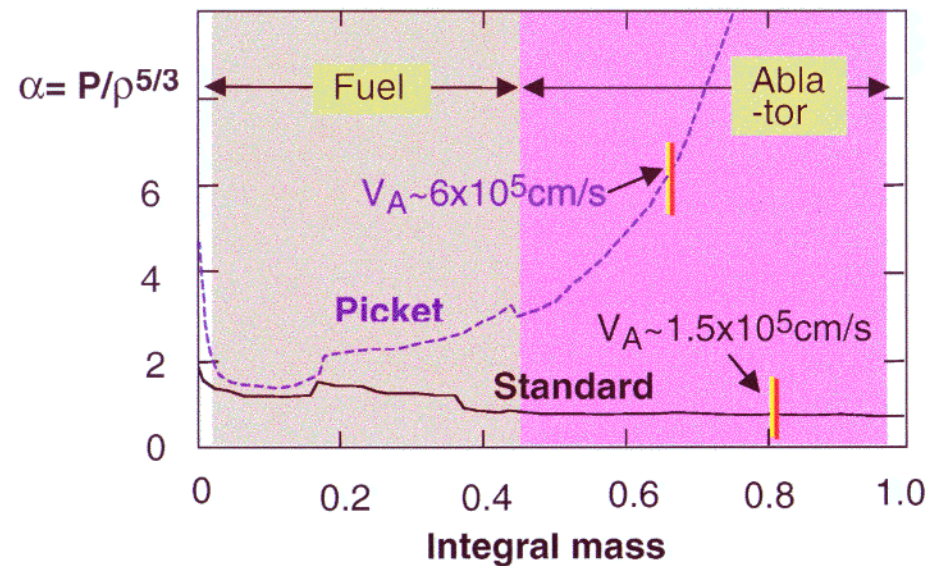
Laser Pulse Shape Power



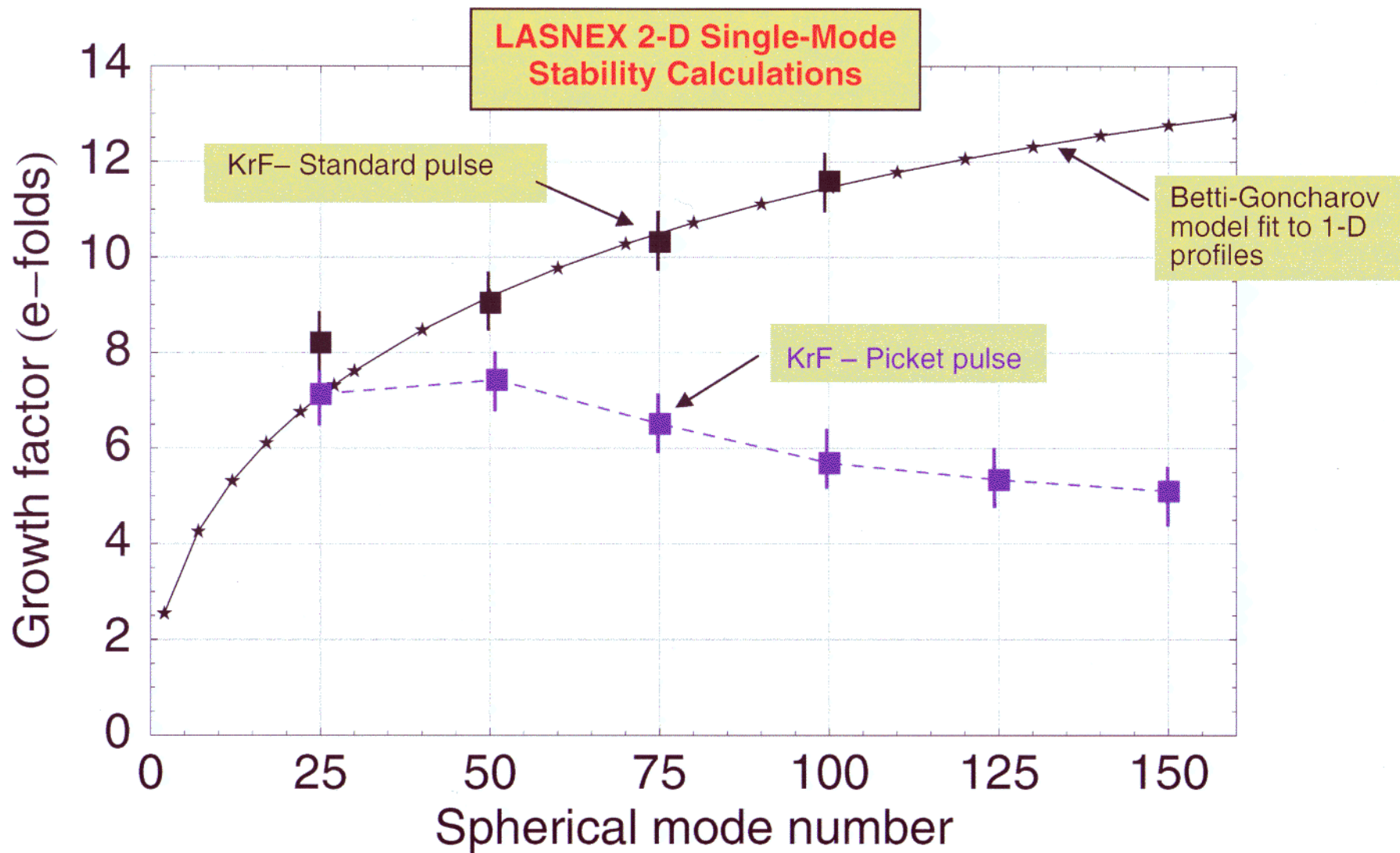
Picket Prepulse:

| Power Fraction | Energy Fraction |
|----------------|-----------------|
| 20% | 2.1% |

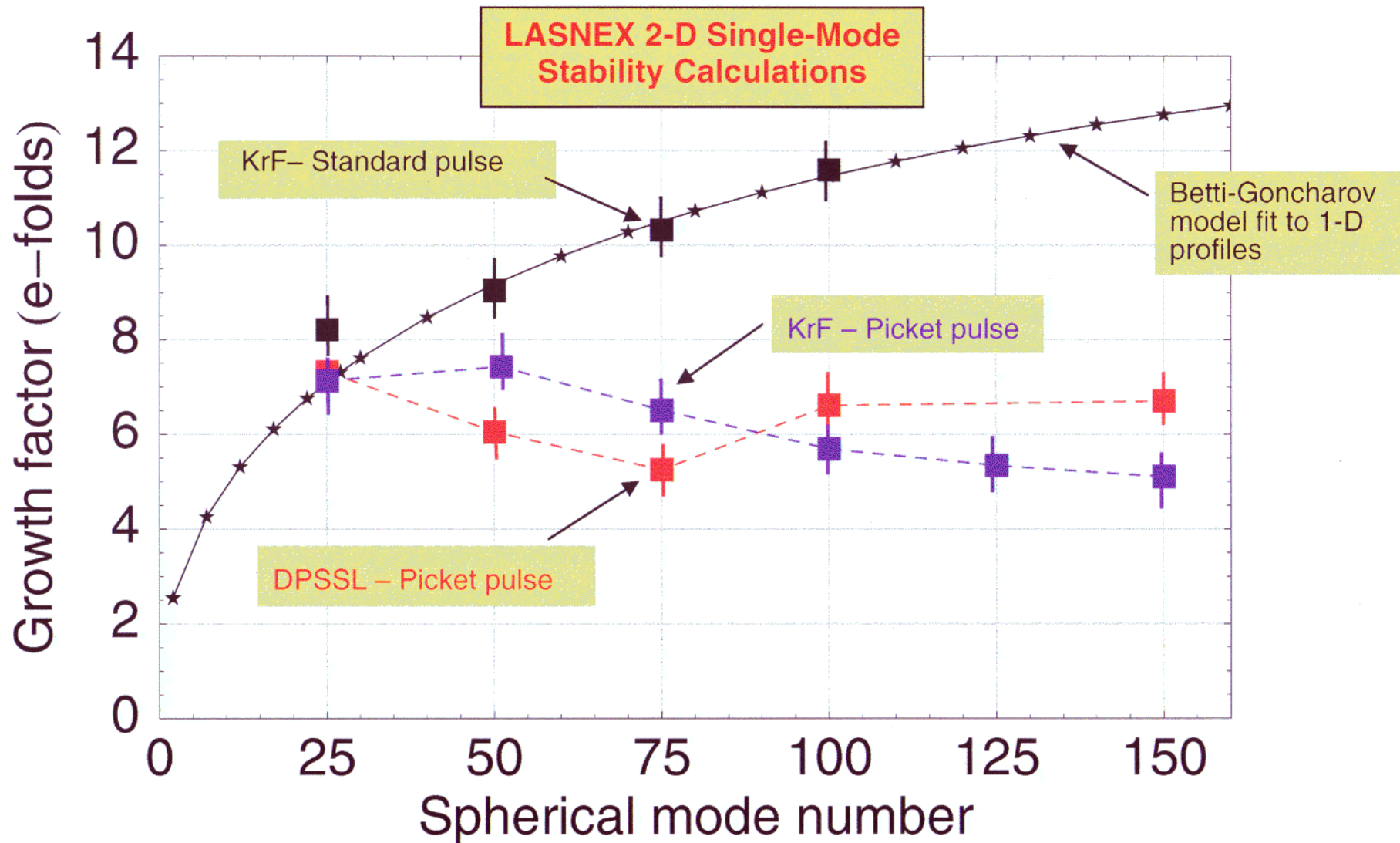
Adiabatic Profile at Time of ~30% Peak Drive Power



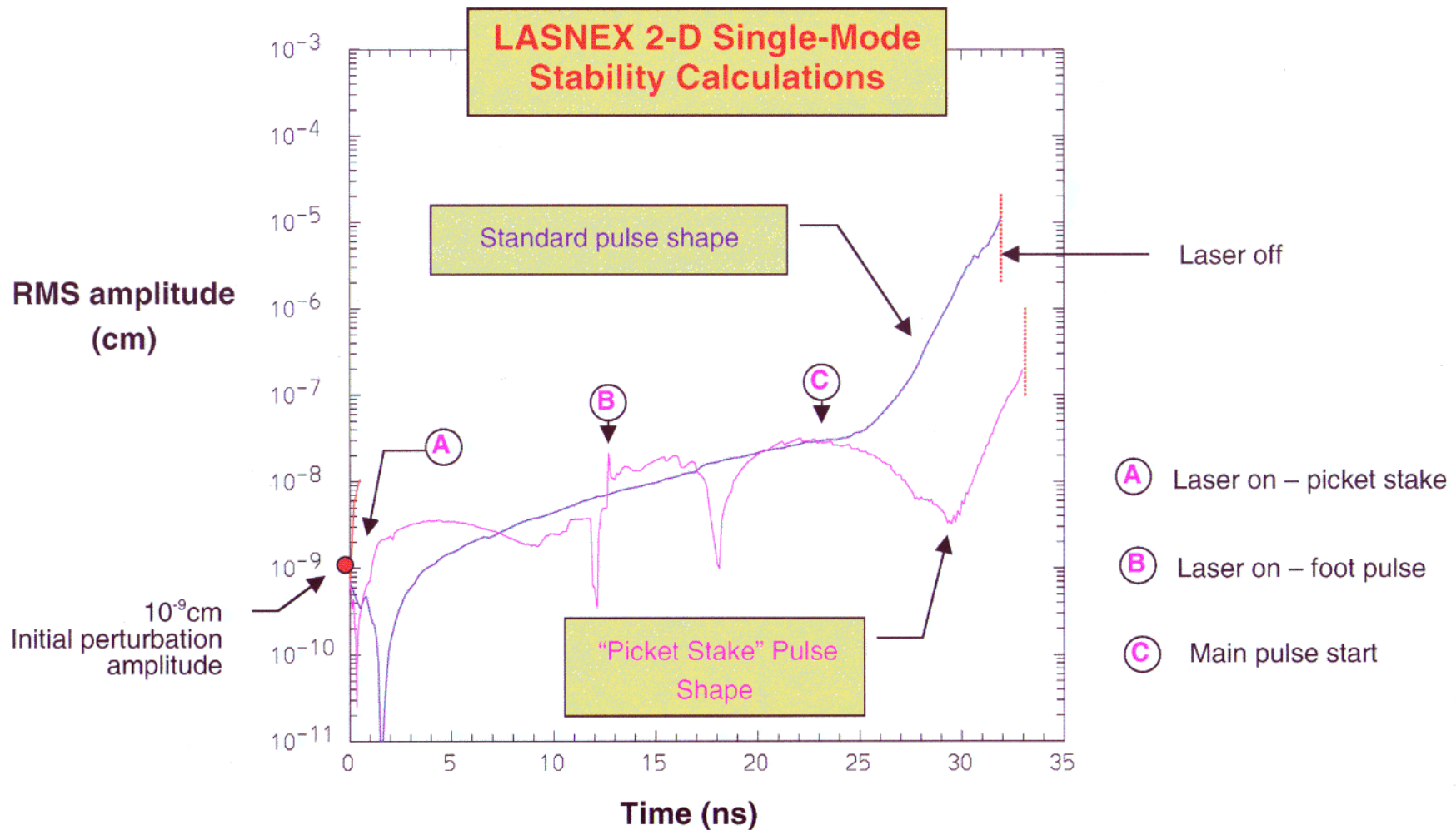
Shaping the KrF Laser Pulse Offers Large Improvements in Stability for the Same Direct Drive Target



KrF and DPSSL Picket Pulses Give Comparable Stability.

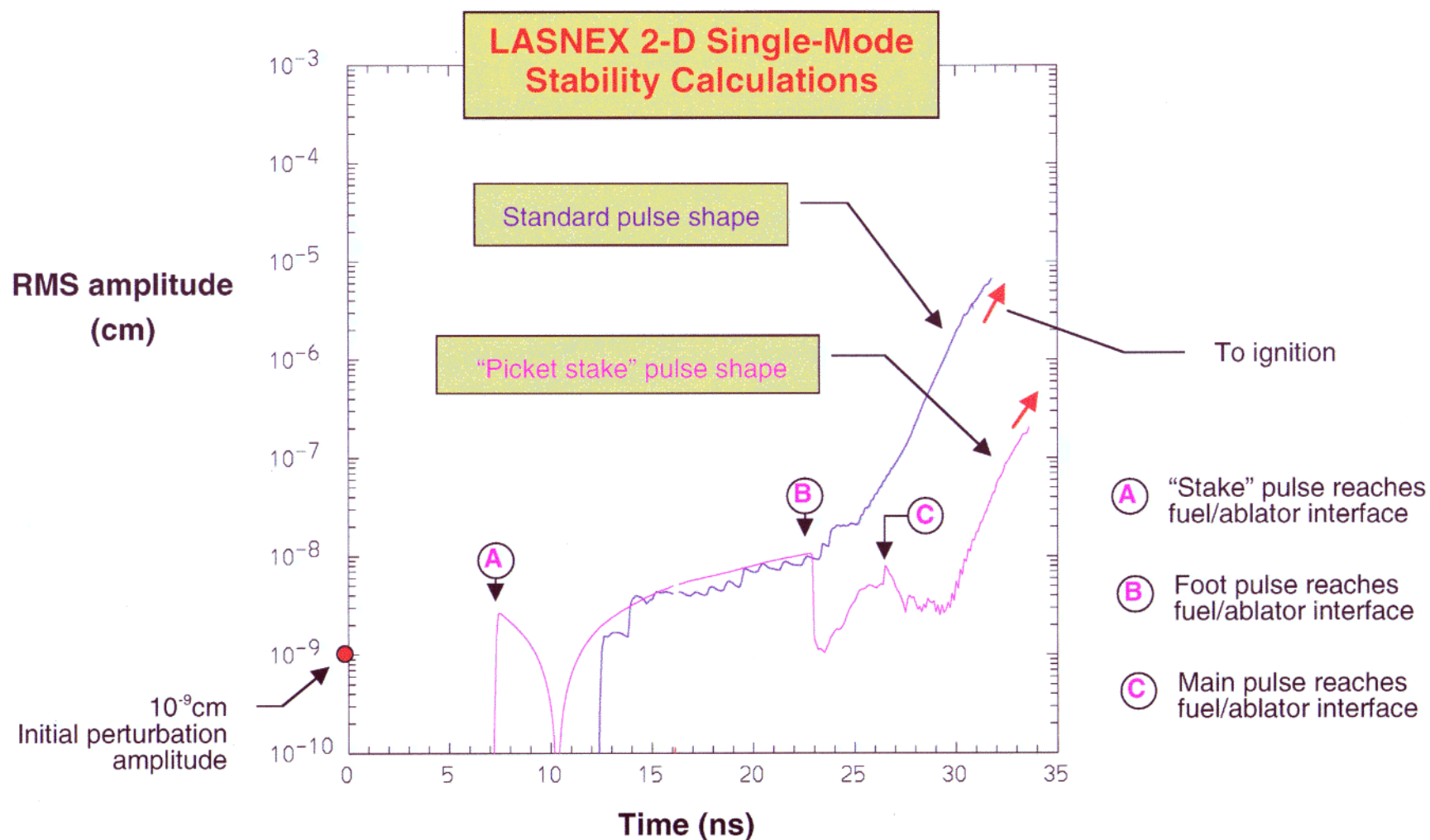


Instability Growth at the Ablation Front @ $\ell = 50$ – Note Multiple Shock Behavior of Amplitude Growth for Picket Pulse



Instability Growth at the Fuel/Abl Interface @ $\ell = 50$

– Note Multiple Shock Behavior of Amplitude Growth for Picket Pulse



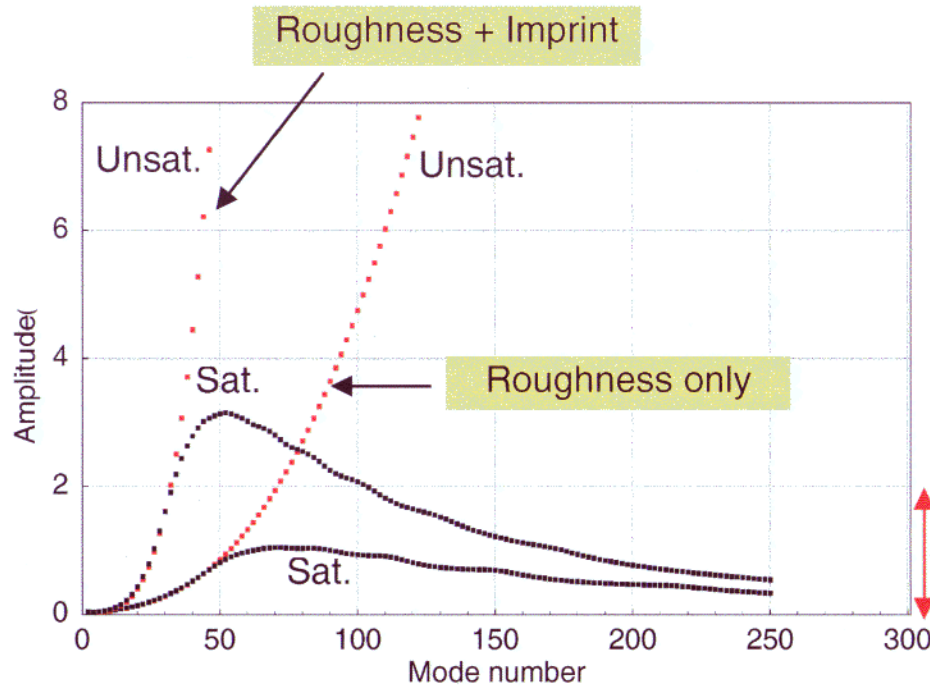
Instability Growths for Tailored Adiabatic Capsules are Less Sensitive to Assumptions of Saturation Models



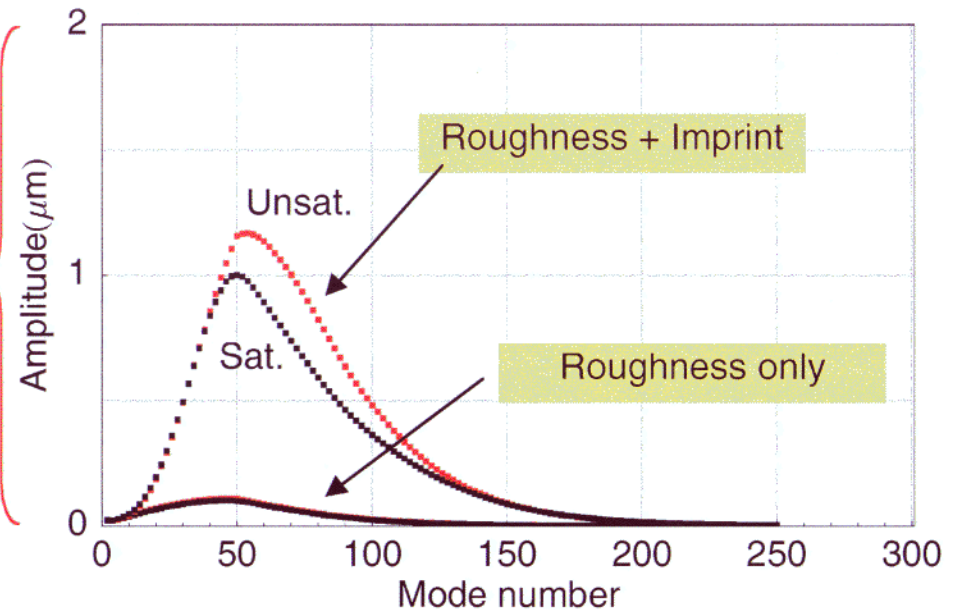
Haan saturation model:

$$a(t_{final}, l) = a_0(l) \underbrace{\text{Exp}\left[\int_{t=0}^{t=t_{sat}} \gamma(t, l) dt\right]}_{\text{Exponential growth to saturation}} \underbrace{\left(1 + \int_{t=t_{sat}}^{t=t_{final}} \gamma(t, l) dt\right)}_{\text{Linear (with time) growth after saturation}}$$

where $t = t_{sat}$ when $a(t_{sat}, l) = C_{Haan} \mathcal{N}(2\pi l)$

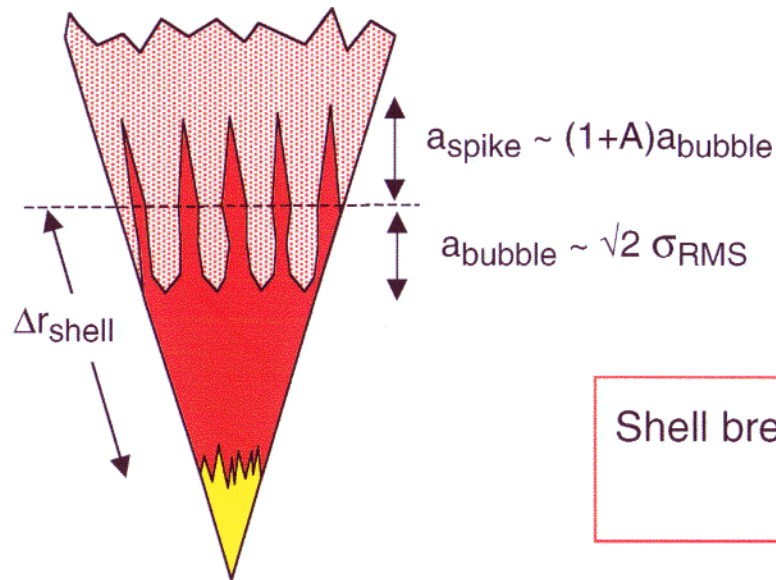


A. Standard Pulse Shape



C. Large Prepulse

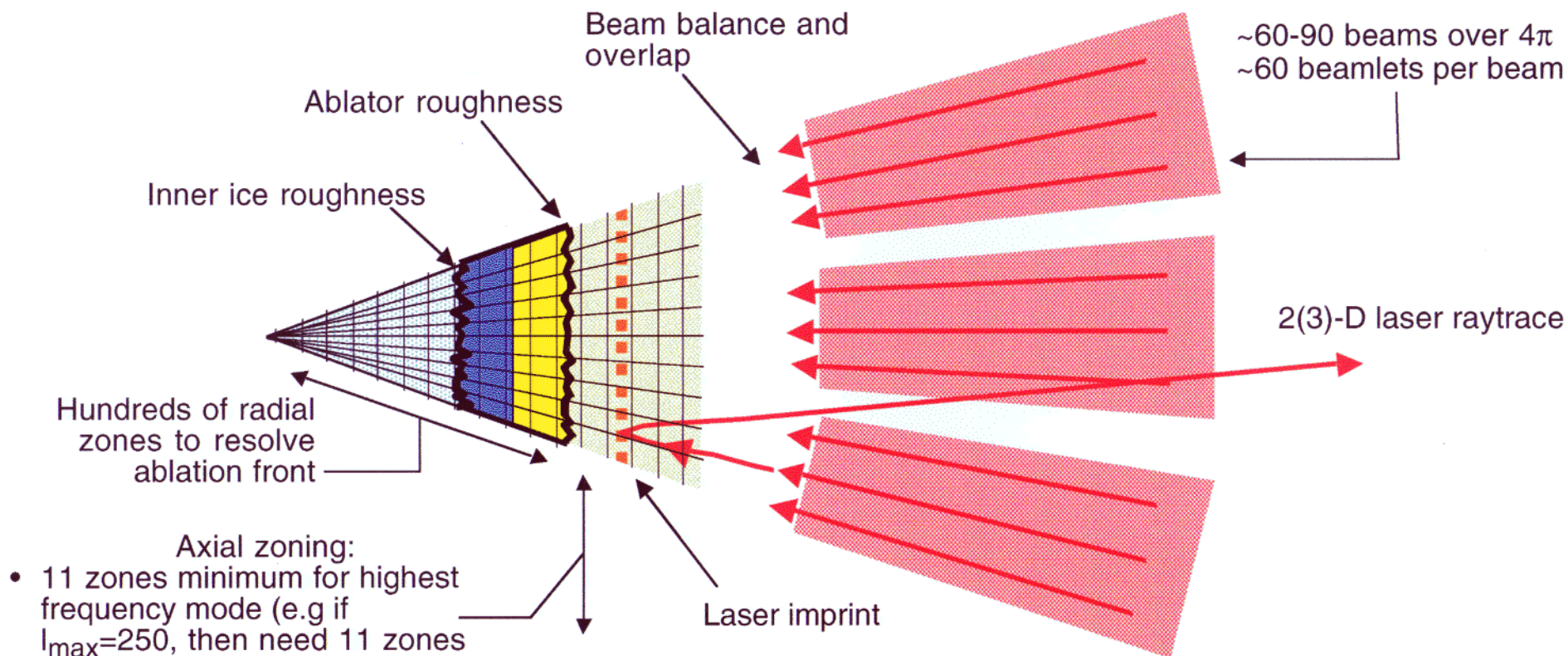
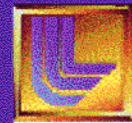
Resulting Shell Breakup Fraction at Late Time is Modest for the Tailored Adiabatic Case



Shell breakup fraction = $a_{\text{bubble}} / \Delta r_{\text{shell}}$
(Require $\leq 30\%$)

| Pulse Shape | Laser (MJ) | Yield (MJ) | Gain | Max Shell Breakup Fraction | |
|------------------------------------|------------|------------|------|----------------------------|---------------------|
| | | | | Roughness Only | Roughness + Imprint |
| A. Standard | 2.4 | 430 | 180 | 0.83 | 1.83 |
| C. Large prepulse (large "picket") | 3.1 | 360 | 110 | 0.015 | 0.15 |

No One has Yet Performed a Full, End-to-End, 2(3)-D Multimode Implosion with Real 2(3)-D Beams



- Axial zoning:
 - 11 zones minimum for highest frequency mode (e.g if $I_{\max}=250$, then need 11 zones over 0.72°)
 - But wedge angle must accommodate $\lambda/2$ for lowest frequency mode (e.g if $I_{\min}=5$, then need 36° wedge angle)

This example would need

- ~550 axial zones,
- ≥ 500 radial zones
- $\geq 275,000$ zones total
- ≥ 5500 rays (≥ 10 per axial zone)
- Weeks of Lasnex run time

